Ma 529 Lecture XIII

Application of the Residue Theorem to the Evaluation of Real Integrals

A.
$$I=\int_0^{2\pi}R(cos\theta,sin\theta)d\theta.$$

Here R(s,t) is a rational function of s and t. We assume that $R(cos\theta, sin\theta)$ is finite on $0 \le \theta \le 2\pi$.

Let $z = e^{i\theta}$. Then the interval $0 \le \theta \le 2\pi$ yields the curve C: |z| = 1.

Now
$$cos\theta = \frac{1}{2}(e^{i\theta} + e^{-i\theta}) = \frac{1}{2}(z + \frac{1}{z})$$
 and $sin\theta = \frac{1}{2i}(z - \frac{1}{z})$; also $dz = ie^{i\theta}d\theta = izd\theta$ or $d\theta = (\frac{1}{i})(\frac{dz}{z})$. Thus

$$I = \oint_C R\left[rac{1}{2}\left(z+rac{1}{z}
ight),\,rac{1}{2i}\left(z-rac{1}{z}
ight)
ight] \;\; ext{where} \; C\colon |\; z\; | \; = 1.$$

I may be evaluated by residues.

Example: Evaluate
$$\int_0^{2\pi} \frac{\cos 2\theta d\theta}{1-2p\cos \theta+p^2}$$
 $-1 .$

Note that by adding and subtracting 2p, the denominator of the integrand can be written in either of two equivalent forms:

$$egin{aligned} 1-2p\,cos heta+p^2&=1-2p+p^2+2p-2p\,cos heta&=(1-p)^2+2p(1-cos heta)\ &=1+2p+p^2-2p\,cos heta&=(1+p)^2-2p(1+cos heta). \end{aligned}$$

From the first of these it is clear that if $0 \le p < 1$, the denominator is different from zero for all values of θ ; and from the second it is clear that if $-1 , the denominator is also different from zero for all values of <math>\theta$. Therefore the integrand is finite on $0 \le \theta \le 2\pi$ for -1 . Now

$$cos2 heta=rac{e^{2i heta}+e^{-2i heta}}{2}=rac{z^2+z^{-2}}{2}.$$

Therefore the integral becomes:

$$\oint_C \frac{z^2 - z^{-2}}{2} \frac{1}{1 - \frac{2p(z + z^{-1})}{2} + p^2} \frac{dz}{iz}$$

$$= \oint_C \frac{z^4 + 1}{2z^2} \frac{z}{z - pz^2 - p + p^2 z} \frac{dz}{iz}$$

$$=\oint_C rac{(1+z^4)dz}{2iz^2(1-pz)(z-p)} \ .$$

There are three poles in the integrand at 0, $\frac{1}{p}$, and p. Now -1

so that $\frac{1}{p}$ is outside of |z| = 1. Therefore we may disregard this pole. For the other two poles we have:

$$\mathop{Res}\limits_{p} = \mathop{lim}\limits_{z o p} (z-p) \; rac{1+z^4}{2iz^2(1-pz)(z-p)} = rac{1+p^4}{2ip^2(1-p^2)} \, .$$

The pole at 0 is second order and therefore

$$\begin{split} Res &= \lim_{z \to 0} \frac{d}{dz} \left[z^2 \, \frac{1 + z^4}{2iz^2(z - pz^2 - p + p^2z)} \right] \\ &= \lim_{z \to 0} \frac{(z - pz^2 - p + p^2z)(4z^3) - (1 + z^4)(1 - 2pz + p^2)}{2i(z - pz^2 - p + p^2z)^2} \\ &= -\left(\frac{1 + p^2}{2ip^2}\right) \\ I &= 2\pi i \, \left[\frac{1 + p^4}{2ip^2(1 - p^2)} - \frac{1 + p^2}{2ip^2} \right] = \frac{2\pi p^2}{1 - p^2} \,. \end{split}$$

B. Improper Integrals of Rational Functions

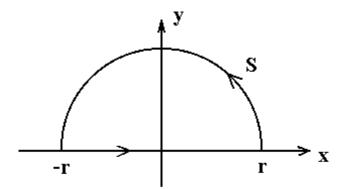
Consider $\int_{-\infty}^{\infty} f(x) dx$ where $f(x) = \frac{P(x)}{Q(x)}$ is a rational function and the degree of $[Q(x)] \geq [P(x)] + 2$. Suppose Q(x) has no real zeros. Such an integral is called an improper integral since the interval of integration is not finite. We define

$$\int_{-\infty}^{\infty} f(x)dx = \lim_{r \to \infty} \int_{-r}^{r} f(x)dx.$$

To evaluate this integral consider the contour integral

$$\oint_C f(z)dz$$

where C is below.



Since f(x) is a rational function, this implies that $f(z) = \frac{P(z)}{Q(z)}$ has finitely many poles in the upper half plane. These are the zeros of Q(z), say $z_1, ..., z_n$. We choose r large enough so that C encloses all of these zeros. Then

$$\oint_C f(z)dz = \int_S f(z)dz + \int_{-r}^r f(x)dx = 2\pi i \sum_{j=1}^n \underset{z_j}{\operatorname{Res}} f(z).$$

Therefore

$$\int_{-r}^{r}f(x)dx=2\pi i\,\sum_{j=1}^{n}\mathop{Res}\limits_{{\mathcal Z}_{j}}f(z)-\int_{S}f(z)dz.$$

It turns out that $\int_S f(z)dz \to 0$ as $r \to \infty$. On S, if we let $z = re^{i\theta}$, then S is represented by r = constant, $0 \le \theta \le \pi$. Since $deg[Q(z)] \ge deg[P(z)] + 2$. Hence

$$\mid f(z) \mid < rac{k}{|z|^2} ext{ for } \mid z \mid = r > r_0,$$

where r_0 and k are sufficiently large. Hence

$$\Big|\int_S f(z)dz\Big| < rac{k}{r^2}\pi r = rac{k\pi}{r} \;\; ext{for}\; r > r_0$$

and this \rightarrow 0 as $r \rightarrow \infty$, so that

$$\int_{-\infty}^{\infty}\!f(x)dx = 2\pi i\,\sum_{j=1}^n \mathop{Res}\limits_{{\mathcal Z}_j} f(z)\,.$$

Example: $\int_{-\infty}^{\infty} \frac{x^2}{(x^2+a^2)(x^2+b^2)} dx$ where a, b > 0.

The poles of $\frac{z^2}{(z^2+a^2)(z^2+b^2)}$ are at $z=\pm ai$, $\pm bi$. Of these only z=ai and z=bi lie in the upper half plane. At z=ai the residue is:

$$\lim_{z o ai}{(z-ai)}\;rac{z^2}{(z-ai)(z+ai)(z^2+b^2)}=rac{-a^2}{2ai(-a^2+b^2)}=rac{a}{2i(a^2-b^2)}\;.$$

Similarly, at z = bi we have $\frac{b}{2i(b^2 - a^2)}$.

Hence the value of the integral is: $2\pi i \left[\frac{a}{2i(a^2-b^2)} + \frac{b}{2i(b^2-a^2)}\right] = \frac{\pi}{a+b}$.

C.
$$\int_{-\infty}^{\infty} f(x) \cos ax \, dx$$
 and $\int_{-\infty}^{\infty} f(x) \sin ax \, dx$. Here f is as in B.

Note that

$$\int_{-\infty}^{\infty} f(x) cos \ ax \ dx + i \int_{-\infty}^{\infty} f(x) sin \ ax \ dx = \int_{-\infty}^{\infty} f(x) e^{iax} dx.$$

The $\int_{-\infty}^{\infty} f(x)e^{iax}dx$ may be treated as above. Thus

$$\int_{-\infty}^{\infty} f(x)e^{iax}dx = 2\pi i \sum$$
Residues of $e^{iax}f(x)$ at its poles in the upper half plane.

Therefore

$$\int_{-\infty}^{\infty} f(x) cos \ ax \ dx = -2\pi \ \sum$$
 imaginary parts of residues of $e^{iax} f(z)$ at its poles in upper half plane.

Similarly

$$\int_{-\infty}^{\infty} f(x) \sin ax \, dx$$
= $2\pi \sum$ real parts of the residues of $e^{iax} f(z)$ at its poles in upper half plane.

Example.
$$\int_{-\infty}^{\infty} \frac{\cos ax dx}{k^2 + x^2} = \frac{\pi}{k} e^{-ak} \; ; \; \int_{-\infty}^{\infty} \frac{\sin ax dx}{k^2 + x^2} = 0 \quad a > 0, \; k > 0.$$

Now $\frac{e^{iax}}{k^2+z^2}$ has only one pole in the upper half plane, namely a single pole at z=ik. Therefore

$$\mathop{Res}\limits_{z\,=\,ik}\,rac{e^{iaz}}{k^2+z^2}=\!\mathop{lim}\limits_{z
ightarrow ik}(z-ik)\,\,rac{e^{iaz}}{(z-ik)(z+ik)}=rac{e^{-ka}}{2ik}.$$

Therefore

$$\int_{-\infty}^{\infty} rac{e^{iax}}{k^2+x^2} dx = 2\pi i \; rac{e^{-ka}}{2ik} = rac{\pi}{k} \; e^{-ka}$$

and this yields the result.

Laplace Transforms

The Laplace Transform of a (real) function f(t) is given by

$$\mathcal{L}[f(t)] = \phi(s) = \int_0^\infty e^{-st} f(t) dt$$

for all s such that this integral converges.

Example. $f(t) = e^{at}$

$$\begin{split} \mathcal{L}[f(t)] &= \mathcal{L}[e^{at}] = \int_0^\infty e^{-st} e^{at} \ dt = \int_0^\infty e^{-(s-a)^t} dt \\ &= \lim_{R \to \infty} \int_0^R e^{-(s-a)t} \ dt = \lim_{R \to \infty} \left. \frac{e^{-(s-a)t}}{a-s} \right|_0^R \\ &= \lim_{R \to \infty} \left. \frac{e^{-(s-a)R} - 1}{a-s} = \frac{1}{s-a} \right. \text{ if } s > a. \end{split}$$

Therefore $\mathcal{L}[e^{at}] = \frac{1}{s-a}$.

We may define an inverse operator \mathcal{L}^{-1} where $\mathcal{L}^{-1}[\phi(s)] = f(t) \Leftrightarrow \mathcal{L}[f(t)] = \phi(s)$. **Thus**

$$\mathcal{L}^{-1}igl[rac{1}{s-a}igr] = e^{at}.$$

Laplace transforms are very useful in solving ordinary and partial differential equations. To illustrate the power of the method, consider the problem: $\frac{dy}{dt} + y = e^{2t}$; y(0) = -1.

$$rac{dy}{dt} + y = e^{2t}; \ y(0) = -1.$$

We take Laplace transforms of both sides of the equation to get

$$\mathcal{L}ig[rac{dy}{dt}ig] + \mathcal{L}[y] = \mathcal{L}[e^{2t}] = rac{1}{s-2}$$

Now
$$\mathcal{L}\left\{\frac{dy}{dt}\right\} = \int_0^\infty \frac{dy}{dt} e^{-st} dt = y e^{-st} \Big|_0^\infty + s \int_0^\infty y e^{-st} dt$$
.
Let $dv = \frac{dy}{dt} dt$ and $u = e^{-st}$ so that $du = -s e^{-st} dt$ and $v = y$. Then

 $\mathcal{L}[y'] = -y(0) + s\mathcal{L}[y]$. The D.E. then leads to

$$(s+1)\mathcal{L}[y] + 1 = rac{1}{s-2} \ rac{1}{s-2} - 1 = rac{-s+3}{s-2}$$

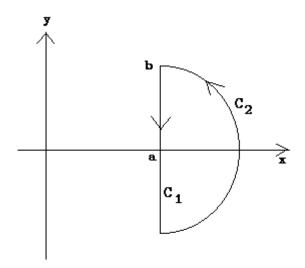
or
$$\mathcal{L}[y] = \frac{-s+3}{(s-2)(s+2)}$$

$$\frac{-s+3}{(s-2)(s+1)} = \frac{A}{s-2} + \frac{B}{s+1} \quad s = 2 \ \to \ \frac{1}{3} = A, \quad s = -1 \ \to \ \frac{4}{-3} = B$$

$$L[y] = \frac{\frac{1}{3}}{s-2} + \frac{\frac{-4}{3}}{s+1} \ \to \quad y = \frac{1}{3}e^{2t} - \frac{4}{3}e^{-t} \\ y' = \frac{2}{3}e^{2t} + \frac{4}{3}e^{-t} \\ \end{pmatrix} y' + y = e^{2t}.$$

This method may be extended to higher order equations and to partial differential equations. Note that in order to get y(t) we had to find $\mathcal{L}^{-1}[\phi(s)]$. There exists a general approach to finding $\mathcal{L}^{-1}[\phi(s)]$ which is based on the complex variable theory we have developed.

Let $\phi(z)$ be a function of z analytic on the line x=a and in the entire half plane R to the *right* of this line. Moreover, let $|\phi(z)| \to 0$ uniformly as z becomes infinite through this plane. Let s be any point in the half plane R. We can choose a semicircular contour $C = C_1 \bigcup C_2$ as in the figure below and apply the Cauchy Integral Formula.



$$\phi(s) = \frac{1}{2\pi i} \oint_C \frac{\phi(z)}{z-s} dz = \frac{1}{2\pi i} \int_{a+ib}^{a-ib} \frac{\phi(z)}{z-s} dz + \frac{1}{2\pi i} \int_{C_2} \frac{\phi(z)}{z-s} dz.$$
 It can be shown that
$$\lim_{b\to\infty} \frac{1}{2\pi i} \int_{C_2} \frac{\phi(z)}{z-s} dz = 0 \text{ so that}$$

$$\phi(s)=rac{1}{2\pi i}{\int_{a-i\infty}^{a+i\infty}rac{\phi(z)}{s-z}dz}$$

Now we shall try to determine the function f(t) whose Laplace Transform is $\phi(s)$.

$$\mathcal{L}^{-1}[\phi(s)] = f(t) = \mathcal{L}^{-1}igg\{rac{1}{2\pi i} \!\int_{a-i\infty}^{a+i\infty} rac{\phi(z)}{s-z} dzigg\}$$

Hence

$$egin{aligned} f(t) &= rac{1}{2\pi i} \int_{a-i\infty}^{a+i\infty} \mathcal{L}^{-1} \Big\{ rac{\phi(z)}{s-z} \Big\} dz \ &= rac{1}{2\pi i} \int_{a-i\infty}^{a+i\infty} \phi(z) \mathcal{L}^{-1} \Big\{ rac{1}{s-z} \Big\} dz \ &= rac{1}{2\pi i} \int_{a+i\infty}^{a+i\infty} \phi(z) e^{zt} dz. \end{aligned}$$

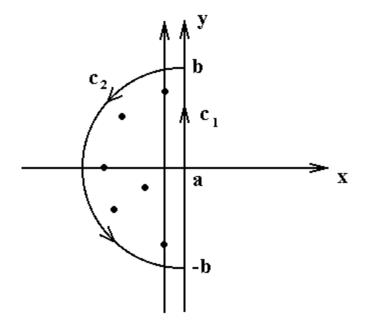
since \mathcal{L}^{-1} only operates on the variable s.

Thus f(t) is given by a line integral in the complex plane taken along a vertical line to the right of all singularities of the transform $\phi(s)$. This complex line integral may be evaluated by the method of residues. A semicircular contour whose diameter is the segment from a-ib to a+ib and whose radius b is large is used. Thus

Theorem: If the Laplace Transform $\phi(s)$ is an analytic function of s except at a finite number of poles each of which lies to the left of the vertical line $Re\ s=a$ and if $s\phi(s)$ is bounded as s becomes infinite through the half plane $Re\ z\leq a$, then

$$\mathcal{L}^{-1}\{\phi(s)\}=\sum$$
 residues of $\phi(s)e^{st}$ at each of its poles.

Remark. Use



Example: What is $\mathcal{L}^{-1}\Big[rac{1}{(s+a)^2+b^2}\Big]$?

Using the theorem we need only compute the residues of $\frac{e^{st}}{(s+a)^2+b^2}$. There are only two first order poles at $-a \pm ib$. At s = -a + ib we have \lim

first order poles at
$$-a\pm ib$$
. At $s=-a+ib$ we have $\lim_{s\to -a+bi} \frac{[s-(-a+ib)]e^{st}}{[s-(-a-ib)][s-(-a-ib)]} = \frac{e^{(-a+ib)t}}{2ib}$.

At s = -a - ib we have:

$$\lim_{s \to -a-bi} \ \frac{[s-(-a-ib)]e^{st}}{[s-(-a+ib)][s-(-a-ib]} = \frac{e^{(-a-ib)t}}{-2ib} \ .$$

Therefore

$$egin{align} f(t) = \mathcal{L}^{-1}\{\phi(s)\} &= rac{e^{(-a+ib)t}}{2ib} + rac{e^{(-a-ib)t}}{-2ib} = e^{-at}\left\{rac{e^{ibt}-e^{-ibt}}{2ib}
ight\} \ &= rac{e^{-at}sinbt}{b} \,. \end{split}$$